# Numerical Simulation of Semi-Elliptical Axial Crack in Pipe Bend Using XFEM

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#### ABSTRACT

In this work, XFEM is employed to obtain the stress intensity factors (SIFs) of a semi elliptical part through thickness axial crack. In XFEM, additional functions are employed to enrich the displacement approximation using partition of unity approach. Level set functions are approximated using higher order shape functions in the crack front elements to ensure the accurate modeling of the crack. The axial crack is placed either on the inner or the outer surface in an internally pressurized pipe bend. The SIFs are extracted from XFEM solution by domain type interaction integral approach for a wide range of geometry parameters like bend radius ratio, cross sectional radius ratio and relative crack length. The results obtained by XFEM approach are compared with those obtained by FEM. These simulations show that the orientation and type of crack in pipe bend has a significant effect on the SIF.

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# **1 INTRODUCTION**

**C**RACKS are found in all machine components at micro or macro level. The crack/defect assessment methods are needed for the safe operation of plant components. Pipe bend is important part of pipeline systems in the chemical and nuclear industry. It has more geometric variables involved than straight pipe. These are generally known to be the most economical means of changing directions while providing flexibility, and end reactions to systems within the allowable limits.

A number of analyses exist for straight thick-walled cylinders with semi-elliptical cracks under internal pressure [1-7]. Thick-walled curved tubes and elbows are also presented but the studies on cracked elbow and tubing are very limited. Chattopadhyay et al. [8-9], Yahiaoui et al. [10], Saxena and Ramakrishnan [11] performed the studies on through thickness circumferential and longitudinal cracks under bending loads. Only few standard shape crack problems have the closed form analytical solutions. Therefore, the numerical methods are the only choice left to analyze real life crack problems. Although, numerical methods can simulate these problems but the accurate and efficient simulation of these cracks still remains a most challenging tasks for the fracture mechanics community. It is known that during fatigue loading, an initial crack of semi-elliptical shape grows up to the critical size. Hence, the analysis of semi-elliptical crack in pipe is quite important from the failure point of view. Therefore, the semi-elliptical crack placed either on the outer or the inner surface of the pipe bend has been considered in the present work.



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To simulate cracked structures/components, a number of methods such as finite element method (FEM) [12-16], boundary element method (BEM) [17-18], and meshfree methods [19-20] have been developed over the years. In FEM, the geometry is normally modeled by an adequate mesh as it requires that a crack must coincide with the edges of the finite elements i.e. a conformal mesh is required apart from the requirement of special elements to handle crack tip asymptotic stress fields. The finite element simulation of a part through propagating crack in pipe and pipe bend becomes even more difficult, time consuming and inaccurate due to conformal mesh requirement and mapping the variables from previous step to next step. To get rid of these issues, extended finite element method (XFEM) has been developed as an alternative to solve the problems of 3-D cracks. In XFEM, these discontinuities are modeled by adding enrichment functions in the finite element displacement approximation through partition of unity [21-23]. These enrichment functions are obtained from the theoretical background of the problem. Many enrichment functions have been proposed by the researchers [24-25] to model the cracks. Primarily two enrichment functions are required to model a crack in XFEM provides the precise modeling of the cracks. Sukumar *et al.* [26] presented 3-D XFEM formulations for the modeling of mode-I crack problems, and showed that the XFEM results are compared well with the benchmark solutions.

In this work, not only the XFEM has been extended further for the simulation of 3-D cracks in pipe bend but also a non-iterative solution scheme has been adopted to define the geometry of the crack surface. A vector level set approach [27-28] is utilized for analyzing the 3-D cracks. Domain based interaction integral approach is adopted for mixed mode stress intensity factors [29]. The XFEM results are obtained for a wide range of pipe bend parameters like bend radius ratio, radius ratio and relative crack length.

## 2 XFEM METHODOLOGY

Three-dimensional XFEM formulation for a crack is described in the following sub-sections:

#### 2.1 Governing equations

The governing equations of elasto-statics with internal discontinuities are briefly reviewed and an associated weak form is given [23]. Consider a domain, bounded by the boundary is partitioned into three sets:  $\Gamma_u$ ,  $\Gamma_t$  and  $\Gamma_c$  as shown in Fig. 1. Displacements are prescribed on  $\Gamma_u$ , tractions are prescribed on  $\Gamma_t$ . All  $\Gamma_c$  surfaces are assumed to be traction free. The equilibrium conditions and boundary conditions for this problem are given as:

$$\nabla . \sigma + b = 0 \quad \text{in} \quad \Omega \tag{1}$$

$$\sigma \cdot \hat{n} = \bar{t} \quad \text{on } \Gamma_t \tag{1a}$$

$$\sigma \cdot \hat{n} = 0 \quad \text{on} \quad \Gamma_c \tag{1b}$$

where,  $\sigma$  is the Cauchy stress tensor, *u* is the displacement field, *b* is the body force per unit volume and  $\hat{n}$  is the unit outward normal. It is assumed that displacements remain small and the kinematics equations consist of the strain-displacement relation:

$$\varepsilon = \varepsilon(u) = \nabla_s u \tag{2}$$

where,  $\nabla_s$  is the symmetric part of the gradient operator. The boundary conditions

$$u = \overline{u}$$
 on  $\Gamma_{u}$  (3)

The constitutive relation for the elastic material under consideration is given by Hook's law

$$\sigma = D \varepsilon \tag{4}$$

where, *D* is the Hooke's tensor.



# 2.2 XFEM formulation

The XFEM eliminates the burden associated with the conformal mesh requirement for the problems involving cracks. In XFEM, domain is divided into two distinct parts i.e. enriching the finite element approximation by additional functions to model the internal boundaries and mesh generation for the domain.

According to the concept of partition of unity, the standard approximation is enriched with additional functions. The enriched displacement approximation can be written in general form for 3-D crack modeling [32]:

$$\mathbf{u}^{h}(\mathbf{x}) = \sum_{j=1}^{n} N_{j}(\mathbf{x}) \left[ \mathbf{u}_{j} + \underbrace{\left( H(\mathbf{x}) - H(\mathbf{x}_{j}) \right) \mathbf{a}_{j}}_{j \in n_{r}} + \underbrace{\sum_{\alpha=1}^{4} \left( \phi_{\alpha}(\mathbf{x}) - \phi_{\alpha}(\mathbf{x}_{j}) \right) \mathbf{b}_{j}^{\alpha}}_{j \in n_{A}} \right]$$
(5)

where, *n* is the set nodes in the mesh,  $n_r$  is the set of those nodes whose shape function is completely cut by the crack surface, and  $n_A$  is the set of those nodes whose shape function support is partially cut by the crack front,  $u_j$  is the nodal displacement vector associated with the continuous part of the finite element solution,  $a_j$  is the additional degree of freedom vector associated with the Heaviside function, H(X),  $b_j^{\alpha}$  is the additional degree of freedom vector associated with the Heaviside function, H(X),  $b_j^{\alpha}$  is the additional degree of freedom vector associated with asymptotic functions,  $\phi_{\alpha}(X)$ ; H(X) is the Heaviside function defined for those elements which are completely cut by the crack surface, and  $\phi_{\alpha}(X)$  are the asymptotic functions defined for those elements which are partially cut by the crack front. To model the radial as well as the angular behavior of asymptotic crack-tip stress fields, the following four enrichment functions are used [23]:

$$\phi_{\alpha}(\mathbf{x}) = \{\phi_1, \phi_2, \phi_3, \phi_4\} = \left[\sqrt{r}\cos\frac{\theta}{2}, \sqrt{r}\sin\frac{\theta}{2}, \sqrt{r}\cos\frac{\theta}{2}\sin\theta, \sqrt{r}\sin\frac{\theta}{2}\sin\theta\right]$$
(6)

The potential energy function for a solid mechanics problem can be written as:

$$\Pi(\mathbf{u}) = \frac{1}{2} \int_{\Omega} \boldsymbol{\varepsilon}(\mathbf{u}) : \mathbf{C} : \boldsymbol{\varepsilon}(\mathbf{u}) d\Omega - \int_{\Omega} \mathbf{b} \cdot \mathbf{u} d\Omega - \int_{\Gamma_t} \overline{\mathbf{t}} \cdot \mathbf{u} d\Gamma$$
(7)

By substituting the trial and test functions from Eq. (5) in the above equation and using the arbitrariness of the nodal variables, the following set of discrete equations are obtained

$$\begin{bmatrix} \tilde{\mathbf{K}} \end{bmatrix} \{ \tilde{\mathbf{d}} \} = \{ \tilde{\mathbf{f}} \}$$
<sup>(8)</sup>

where,  $\tilde{\mathbf{K}}$  is the global stiffness matrix,  $\tilde{\mathbf{d}}$  is the vector of nodal unknown variables, and  $\tilde{\mathbf{f}}$  is external force vector. The elemental  $\tilde{\mathbf{K}}$  and  $\tilde{\mathbf{f}}$  vectors are given as:

$$\tilde{\mathbf{f}}^{\epsilon} = \left\{ \tilde{\mathbf{f}}^{a}_{i} \quad \tilde{\mathbf{f}}^{b}_{i} \quad \tilde{\mathbf{f}}^{b1}_{i} \quad \tilde{\mathbf{f}}^{b2}_{i} \quad \tilde{\mathbf{f}}^{b3}_{i} \quad \tilde{\mathbf{f}}^{b4}_{i} \right\}^{T}$$
(9)

$$\tilde{\mathbf{K}}_{ij}^{e} = \begin{bmatrix} \tilde{\mathbf{K}}_{ij}^{uu} & \tilde{\mathbf{K}}_{ij}^{ua} & \tilde{\mathbf{K}}_{ij}^{ub} \\ \tilde{\mathbf{K}}_{ij}^{au} & \tilde{\mathbf{K}}_{ij}^{au} & \tilde{\mathbf{K}}_{ij}^{ab} \\ \tilde{\mathbf{K}}_{ij}^{bu} & \tilde{\mathbf{K}}_{ij}^{bu} & \tilde{\mathbf{K}}_{ij}^{bb} \end{bmatrix}$$
(10)

where,

$$\tilde{\mathbf{K}}_{ij}^{\alpha\beta} = \int_{\Omega^{e}} \left( \mathbf{B}_{i}^{\alpha} \right)^{T} \mathbf{C} \mathbf{B}_{j}^{\beta} d\Omega \quad \left( \alpha, \beta = u, a, b \right)$$
(11)

$$\tilde{\mathbf{f}}_{i}^{u} = \int_{\Gamma_{i}} N_{i} \,\overline{\mathbf{t}} \, d\Gamma + \int_{\Omega^{e}} N_{i} \, \mathbf{b} \, d\Omega \tag{12}$$

$$\tilde{\mathbf{f}}_{i}^{a} = \int_{\Gamma_{i}} N_{i} \left( H\left(\mathbf{x}\right) - H\left(\mathbf{x}_{i}\right) \right) \overline{\mathbf{t}} d\Gamma + \int_{\Omega^{c}} N_{i} \left( H\left(\mathbf{x}\right) - H\left(\mathbf{x}_{i}\right) \right) \mathbf{b} d\Omega$$
(13)

$$\tilde{\mathbf{f}}_{i}^{ba} = \int_{\Gamma_{i}} N_{i} \left( \phi_{\alpha}(\mathbf{x}) - \phi_{\alpha}(\mathbf{x}_{i}) \right) \overline{\mathbf{t}} d\Gamma + \int_{\Omega^{e}} N_{i} \left( \phi_{\alpha}(\mathbf{x}) - \phi_{\alpha}(\mathbf{x}_{i}) \right) \mathbf{b} d\Omega$$
(14)

Using shifted enrichment [32],  $\mathbf{B}_{i}^{u}$ ,  $\mathbf{B}_{i}^{a}$  and  $\mathbf{B}_{i}^{b}$  are defined as:

$$\mathbf{B}_{i}^{u} = \begin{bmatrix} N_{i,x} & 0 & 0 \\ 0 & N_{i,y} & 0 \\ 0 & 0 & N_{i,z} \\ 0 & N_{i,z} & N_{i,y} \\ N_{i,z} & 0 & N_{i,x} \\ N_{i,y} & N_{i,x} & 0 \end{bmatrix}$$
(15)

$$\mathbf{B}_{i}^{a} = \begin{bmatrix} \left(N_{i}\left(H(\mathbf{x})-H(\mathbf{x}_{i})\right)\right)_{,x} & 0 & 0\\ 0 & \left(N_{i}\left(H(\mathbf{x})-H(\mathbf{x}_{i})\right)\right)_{,y} & 0\\ 0 & 0 & \left(N_{i}\left(H(\mathbf{x})-H(\mathbf{x}_{i})\right)\right)_{,z}\\ 0 & \left(N_{i}\left(H(\mathbf{x})-H(\mathbf{x}_{i})\right)\right)_{,z} & \left(N_{i}\left(H(\mathbf{x})-H(\mathbf{x}_{i})\right)\right)_{,y}\\ \left(N_{i}\left(H(\mathbf{x})-H(\mathbf{x}_{i})\right)\right)_{,z} & 0 & \left(N_{i}\left(H(\mathbf{x})-H(\mathbf{x}_{i})\right)\right)_{,x}\\ \left(N_{i}\left(H(\mathbf{x})-H(\mathbf{x}_{i})\right)\right)_{,y} & \left(N_{i}\left(H(\mathbf{x})-H(\mathbf{x}_{i})\right)\right)_{,x} & 0 \end{bmatrix}$$
(16)

$$\mathbf{B}_{i}^{b} = \begin{bmatrix} \mathbf{B}_{i}^{b1} & \mathbf{B}_{i}^{b2} & \mathbf{B}_{i}^{b3} & \mathbf{B}_{i}^{b4} \end{bmatrix}$$
(17)

$$\mathbf{B}_{i}^{b} = \begin{bmatrix} \left(N_{i}\left(\phi_{\alpha}(\mathbf{x})-\phi_{\alpha}(\mathbf{x}_{i})\right)\right)_{x} & 0 & 0\\ 0 & \left(N_{i}\left(\phi_{\alpha}(\mathbf{x})-\phi_{\alpha}(\mathbf{x}_{i})\right)\right)_{y} & 0\\ 0 & 0 & \left(N_{i}\left(\phi_{\alpha}(\mathbf{x})-\phi_{\alpha}(\mathbf{x}_{i})\right)\right)_{z}\\ 0 & \left(N_{i}\left(\phi_{\alpha}(\mathbf{x})-\phi_{\alpha}(\mathbf{x}_{i})\right)\right)_{z} & \left(N_{i}\left(\phi_{\alpha}(\mathbf{x})-\phi_{\alpha}(\mathbf{x}_{i})\right)\right)_{y}\\ \left(N_{i}\left(\phi_{\alpha}(\mathbf{x})-\phi_{\alpha}(\mathbf{x}_{i})\right)\right)_{z} & 0 & \left(N_{i}\left(\phi_{\alpha}(\mathbf{x})-\phi_{\alpha}(\mathbf{x}_{i})\right)\right)_{x}\\ \left(N_{i}\left(\phi_{\alpha}(\mathbf{x})-\phi_{\alpha}(\mathbf{x}_{i})\right)\right)_{y} & \left(N_{i}\left(\phi_{\alpha}(\mathbf{x})-\phi_{\alpha}(\mathbf{x}_{i})\right)\right)_{x} & 0 \end{bmatrix}$$
(18)

#### 2.3 Vector level sets for a crack

In vector level set, a crack is defined by two vector functions [26, 28, 30]. These functions are calculated at nodes. Finite element interpolation functions are approximated using these functions. Vector level set is evaluated only for those elements which have their mean centre lying in a narrow band of the crack surface for improving the computational efficiency. Many geometric operations are performed on the data for updating the vector level set. As a result, this crack is entirely described by the nodal data without the need of any physical representation. Thus, the computation of the level set becomes more efficient. The crack geometry is approximated using quadratic twenty node hexahedral elements to improve the accuracy, and the problem domain is approximated by using eight node linear solid elements.

Level set function,  $\varphi(\mathbf{x})$  is a signed distance function, which can be computed from  $\mathbf{f}(\mathbf{x})$  by the following formula  $\varphi(\mathbf{x}) = \|\mathbf{f}(\mathbf{x})\| H(\xi_2, \mathbf{f}(\mathbf{x}))$ , where  $\mathbf{f}(\mathbf{x})$  is the vector between a point  $\mathbf{x}$  and its projection on the crack surface  $(\Gamma_{cr})$ , and is directed from  $\mathbf{x}$  to the crack surface, i.e.  $\mathbf{f}(\mathbf{x}) = \mathbf{x}_f - \mathbf{x}$  where  $\mathbf{x}_f$  is the closest point projection of  $\mathbf{x}$  on the crack surface  $(\Gamma_{cr})$  and  $H(\mathbf{x})$  is the Heaviside function which takes +1 above the crack surface and -1 below the crack surface. The second level set function is a signed distance function with respect to the crack front, and is computed using the following formula  $\psi(\mathbf{x}) = \xi_1 \mathbf{f}(\mathbf{x})$ . A crack surface with the finite elements is represented by

$$\varphi_{approx}\left(\mathbf{x}\right) = \sum_{i=1}^{n} N_{i}\left(\mathbf{x}\right)\varphi(\mathbf{x}_{i})$$
(19)

$$\psi_{approx}\left(\mathbf{x}\right) = \sum_{i=1}^{n} N_{i}\left(\mathbf{x}\right)\psi\left(\mathbf{x}_{i}\right)$$
(20)

The criterion for the finding whether an element belongs to crack front or crack surface, is given as: For crack front elements:  $\psi_{approx}(\mathbf{x}) = 0$  and  $\varphi_{approx}(\mathbf{x}) = 0$ 

approx ( )

For crack surface elements:  $\psi_{approx}(\mathbf{x}) < 0$  and  $\varphi_{approx}(\mathbf{x}) = 0$ .

The enrichment procedure in 3-D is more complicated as compared to 2-D but the basic principle remains the same. In general, it is quite tedious to find nearest point on a curve from an arbitrary point. Therefore, an iterative procedure was adopted in literature to define the geometry of the crack surface [28, 30].

# 2.4 Numerical integration

In case of enriched element, numerical integration needs special treatment due to the presence of discontinuity in the integrand. Integration of enriched elements is achieved by dividing them into several tetrahedrons above and below the surface. To evaluate this, intersection point of element edges and crack front with element surfaces is calculated [33]. Integration of enriched elements generally requires higher order Gauss quadrature. In the present work, three, seven and two points Gauss quadrature scheme is used in crack surface, crack front and rest of the elements, respectively.

# 2.5 Interaction integral

Domain based interaction integral approach has been adopted for the evaluation of individual stress intensity factors  $(K_I, K_{II} \text{ and } K_{III})$  under mixed-mode loading. Fig. 2 shows the domain for *J*-integral. Virtual domain extension approach is used for the evaluation of *J*-integral in 3-D [34]. In this approach, volume integral over the domain surrounding the crack front is expressed in terms of a crack front contour integral. Individual SIFs are extracted from XFEM solution by superimposing two states of the stress. One state (auxiliary state) is assumed to be known while other one is the actual state. The SIF of the known auxiliary state is taken as one for the mode which is being evaluated while the SIFs of remaining two modes in auxiliary states are assigned zero values. Thus, the actual SIF of that mode which is assigned one value is in auxiliary state, remains unknown, and can be evaluated by solving an interaction integral [27, 34-37]]. The same procedure is adopted for the other modes also. Now, the path independent *J*-integral for a homogeneous cracked body is given as:

$$J = \int_{\Gamma} \left( W \ \delta_{1j} - \sigma_{ij} \frac{\partial u_i}{\partial x_1} \right) n_j d\Gamma$$
(21)

*J*-integral is not evaluated over the crack surface since it is traction free. Fig. 2 shows the surface contour integral which is virtually extended along the crack front. Green divergence theorem is used to convert a closed surface integral into the volume integral by taking its gradient as discussed by Pathak et al. [30]. *J*-integral is equal to the energy release rate for elastic materials. For a mixed-mode case, it is given as:

$$J = G = \frac{1}{E} \left( K_{I}^{2} + K_{II}^{2} \right) + \frac{1}{2\mu} K_{III}^{2}$$
(22)

where,  $E' = \frac{E}{1-v^2}$  and  $\mu = \frac{E}{2(1+v)}$ 

The stress intensity factors  $K_I$ ,  $K_{II}$  and  $K_{III}$  are computed in the direction of fatigue crack growth. Now consider two independent equilibrium state of a cracked body. State 1 is taken as actual state for the given boundary conditions while state 2 is an auxiliary state. The J -integral for two superposed states is provided as:

$$J^{tot} = \int_{\Gamma} \left[ W^{tot} \,\delta_{1j} - \left(\sigma_{ij}^{(1)} + \sigma_{ij}^{(2)}\right) \frac{\partial \left(u_i^{(1)} + u_i^{(2)}\right)}{\partial x_1} \right] n_j \, d\,\Gamma$$
(23)

where,  $W^{tot} = \frac{1}{2} \left( \sigma_{ij}^{(1)} + \sigma_{ij}^{(2)} \right) \left( \varepsilon_{ij}^{(1)} + \varepsilon_{ij}^{(2)} \right)$ 

Eq. (23) can be expanded, and rearranged to yield

$$J^{tot} = J^{(1)} + J^{(2)} + M^{(1,2)}$$
(24)

In Eq. (24),  $J^{(1)}$ ,  $J^{(2)}$  are the *J*-integral for states 1 and 2 and  $M^{(1,2)}$  is the interaction integral for two equilibrium states.

$$M^{(1,2)}(t) = \frac{2(1-\nu^2)}{E} \left[ K_I K_I^{aux} + K_{II} K_{II}^{aux} \right] + \frac{1}{2\mu} K_{III} K_{III}^{aux}$$
(25)



Fig. 2 Hollow cuboidal domain around penny crack.

# **3 NUMERICAL RESULTS AND DISCUSSION**

In the present work, the stress intensity factors for a cracked pipe bend are evaluated under internal pressure. An axial semi-elliptical crack is placed either at extrados or at the intrados. A schematic of the pipe bend is shown in Fig. 3. Pipe bend has been analyzed for different values of bend radius ratio ( $R^*=R/R_1$ ), cross-sectional radius ratio ( $k=R_2/R_1$ ) and relative crack length (a/t). The extent of the pipe bend is taken as 90 degree. In this study, the following values of various parameters are considered: bend radius ratio,  $R^* = R/R_1 = 5$ , 6.25, 7.5 and 8.75, cross-sectional radius ratio,  $k = R_2/R_1 = 1.25$ , 1.375, 1.5 and 1.625.

The relative crack depth (a/t), where  $t = (R_2 - R_1)$  is the wall thickness, and values of a/t is taken as 0.5 and 0.75. The crack shape aspect ratio is taken as b/a = 1.6. A tangential straight pipe section (with same radius ratio, k) of length  $3R_2$  is added at each end of the elbow/pipe bend along with displacement constraints to reduce the end effects.

A number of cases of pipe bend with an axial elliptical cracks lying at the intrados as well as at the extrados have been taken for the simulation. The values of stress intensity factors are evaluated at the tip of minor axis, and are compared with the finite element solutions. An extrinsic enrichment technique based on partition of unity has been implemented to capture the discontinuity due to the presence of cracks. To model the curved crack front accurately, it has been divided into one hundred splined-curve parts for all the problems under consideration. The values of SIFs are evaluated for each location of the axial elliptical crack by varying the bend radius ratio, radius ratio and relative crack depth. The results obtained by XFEM are compared with FEM solution obtained in ANSYS using a mesh of 5x100x18 nodes.



Fig. 3 Schematic of pipe bend.

### Case 1: Axial part through crack at extrados external

Pipe bend having an axial semi-elliptical crack lying at the outer surface of the bend has been simulated by proposed XFEM approach. The geometry of the pipe bend along with the crack and boundary conditions is shown in Fig. 4. The material of the bend is assumed to be homogeneous and isotropic with  $E = 200\ 000\ \text{N/mm}^2$  and Poisson ratio

0.3. The minor and major axis of semi-elliptical crack are taken as 0.015 m and 0.024 m. The pipe bend is subjected to an internal pressure of 10 MPa. An eight node linear brick element has been used for meshing the geometry. The dimensions of contour integral are taken as Ve = 0.004 m and b = 0.004 m. A regular mesh of 5x100x18 nodes is used in all simulations.

The variation of SIF with pipe bend radius ratio  $(R^*)$  is shown in Fig. 5 for four different values of crosssectional radius ratio (k). The relative crack depth (a/t) and inner radius of the pipe bend  $(R_1)$  are kept constant for these simulations. Figs. 5(a), 5(b), 5(c) and 5(d) show the variation of SIF with  $R^*$  for k values of 1.25, 1.375, 1.5 and 1.625, respectively. From these figures, it can be seen that the SIF values marginally increases with the increase in  $R^*$  from 5.0 to 8.75. Fig. 6 depicts the variation of stress intensity factor with respect to  $R^*$  for four different values of k. The relative crack depth (a/t) is kept constant and is taken as 0.5. The inner radius of the pipe bend  $(R_1)$ is also kept constant at 0.08 m. Figs. 6(a), 6(b), 6(c) and 6(d) present the variation of SIF with  $R^*$  for the k values of 1.25, 1.375, 1.5 and 1.625, respectively. It can be concluded that the SIF values marginally increases with the increase in bend radius ratio from 5.0 to 8.75. It is also observed that the large relative crack depth results in large SIF values, if rests of pipe bend parameters kept constant.

The variation of SIF with k is shown in Fig.7 for four different values of  $R^*$ . The relative crack depth (a/t) is kept constant, and is taken as 0.75. The inner radius of the pipe bend ( $R_1$ ) is kept constant. Figs. 7(a), 7(b), 7(c) and 7(d) show the variation of SIF with k for  $R^*$  equal to 5.0, 6.25, 7.5 and 8.75, respectively. From these figures, it can be seen that the SIF values decreases with the increase in k from 1.25 to 1.625. Fig.8 depicts the variation of stress intensity factor with k for four different values of  $R^*$ . The relative crack depth (a/t) is kept constant, and is taken as 0.5. The inner radius of the pipe bend ( $R_1$ ) is kept constant for these simulations. The variation of SIF with radius ratio is shown in Figs. 8(a), 8(b), 8(c) and 8(d) for  $R^*$  of 5.00, 6.25, 7.5 and 8.75, respectively. It is observed that the SIF values decreases with the increase in  $R^*$  from 1.25 to 1.625. Fig.8 depicts the large relative crack depth results in high SIF values, if the rest of pipe bend parameters are kept constant. From the SIF, it is observed that the numerical solution obtained by XFEM is found quite close to the solution obtained using Ansys (FEM). It is also found that the relatively large crack depth results in large SIFs, if the rest of parameters are kept constant. It is also found that influence of cross sectional radius ratio (k) is more on SIFs in comparison to bend radius ratio ( $R^*$ ).



Fig. 4 Pipe Bend crack at extrados external location.



**Fig. 5** SIF variation with bend radius ratio ( $R^*$ ) for different values of *k* and *a*/*t*=0.75.









**Fig. 7** SIF variation w.r.t. radius ratio for different values of R\* and *a/t*=0.75.







# Case II: Axial part through crack at extrados internal

In the present case, pipe bend dimensions, crack length, loading, boundary conditions, material properties and meshing remain the same as in case-I. An axial semi-elliptical crack lying at inner surface has been simulated by XFEM. The geometry of the pipe bend along with the crack and boundary conditions is shown in Fig.9. The dimensions of interaction integral contour are taken as Ve = 0.004 m and b = 0.004 m. A regular mesh of  $5 \times 100 \times 18$  nodes has been used in these simulations.

The variation of SIF with pipe bend radius ratio ( $R^*$ ) is shown in Fig.10 for four different values of crosssectional radius ratio (k). The relative crack depth (a/t) is kept constant and is taken as 0.75. The inner radius of the pipe bend ( $R_1$ ) is also kept constant at 0.08 m. Figs. 10(a), 10(b), 10(c) and 10(d) show the variation of SIF with  $R^*$ for four k values of 1.25, 1.375, 1.5 and 1.625, respectively. Fig. 11 depicts the variation of stress intensity factor with respect to  $R^*$  for four different values of k. The value of a/t is kept constant and is taken as 0.5. Figs.11(a), 11(b), 11(c) and 11(d) present the variation of SIF with  $R^*$  for four values k of 1.25, 1.375, 1.50 and 1.625, respectively. It can be concluded that the SIF values marginally increases with increase in bend radius ratio from 5.0 to 8.75. The variation of SIF with radius ratio (k) is shown in Fig. 12 for four different values of bend radius ratio ( $R^*$ ). The relative crack depth (a/t) is taken as 0.75 for these simulations. The inner radius of the pipe bend ( $R_1$ ) is kept constant at 0.08 m. Figs. 12(a), 12(b), 12(c) and 12(d) show the variation of SIF with k for  $R^*$  equal to 5.0, 6.25, 7.5 and 8.75, respectively. From these figures, it can be seen that the SIF values decreases with the increase in k from 1.25 to 1.625.

The variation of stress intensity factor with respect to k for four different values of  $R^*$  is shown in Fig.13. The value of a/t is kept constant at 0.5. Figs. 13(a), 13(b), 13(c) and 13(d) present the variation of SIF with k for four values of  $R^*$  equal to 5.00, 6.25, 7.5 and 8.75, respectively. It can be concluded that the SIF values decreases with increase in bend radius ratio from 1.25 to 1.625. From the SIF plots, it is noticed that higher crack depth results in higher the SIFs, if the rest of parameters are kept constant. It was also found that influence of cross sectional radius ratio is more in comparison to bend radius ratio. It is also found that values of SIF are higher for externally exposed crack in comparison to internally exposed crack. Moreover, it is also observed that the numerical results obtained by XFEM are found in good agreement with those obtained by FEM.



**Fig. 9** Pipe Bend crack at extrados internal location.



# **Fig. 10** SIF variation with bend radius ratio ( $R^*$ ) for different values of *k* and *a/t*=0.75.















# Case III: Axial part through crack at intrados external

In the present case, pipe bend dimensions, crack length, loading, boundary conditions, material properties and meshing remain same as in case-I. An axial semi-elliptical crack lying at the outer surface has been simulated by XFEM. The geometry of the pipe bend along with the crack and boundary conditions is shown in Fig.14. The contour integral dimensions are taken as Ve = 0.004 m and b = 0.004 m. A regular mesh of 5x100x18 nodes is used for these simulations.

The variation of SIF with pipe bend radius ratio  $(R^*)$  is shown in Fig.15 for four different values of crosssectional radius ratio (k). The relative crack depth (a/t) is kept constant and is taken as 0.75. The inner radius of the pipe bend  $(R_1)$  is also kept constant at 0.08 m for these simulations. Figs. 15(a), 15(b), 15(c) and 15(d) show the variation of SIF with bend radius ratio for the cross sectional radius ratio of 1.25, 1.375, 1.5 and 1.625, respectively. From these figures, it can be seen that the SIF values marginally decreases with the increase in bend radius ratio. Fig.16 depicts the variation of stress intensity factor with  $R^*$  for four different values of k. The relative crack depth (a/t) is kept constant, and is taken as 0.5. The variation of SIF with  $R^*$  is shown in Figs. 16(a), 16(b), 16(c) and 16(d) for k equal to 1.25, 1.375, 1.5 and 1.625, respectively. It can be concluded that the SIF values marginally decreases with the increase in bend radius ratio from 5.0 to 8.75. The variation of SIF with radius ratio (k) is shown in Fig. 17 for four different values of bend radius ratio  $(R^*)$ . The relative crack depth (a/t) is kept constant, and is taken as 0.75. The inner radius of the pipe bend  $(R_1)$  is kept constant. Figs. 17(a), 17(b), 17(c) and 17(d) show the variation of SIF with k for  $R^*$  equal to 5.00, 6.25, 7.5 and 8.75, respectively. From these figures, it can be seen that the SIF values decreases with the increase in radius ratio.

The variation of stress intensity factor ( $K_1$ ) with radius ratio is shown in Fig. 18 for four different values of bend radius. The relative crack depth (a/t) is kept constant and is taken as 0.5. The inner radius of the pipe bend ( $R_1$ ) is also kept constant for these simulations. Figs. 18(a), 18(b), 18(c) and 18(d) present the variation of SIF with radius ratio for the bend radius ratio of 5.00, 6.25, 7.5 and 8.75, respectively. It is found that the SIF values decreases with the increase in bend radius ratio. From above plots, it is observed that the numerical solution obtained by XFEM approach is quite close to the solution obtained using FEM (Ansys). It is also found that relatively higher crack depth results in higher SIFs if the rest of parameters are kept constant. It was also found that influence of cross sectional radius ratio is more in comparison to bend radius ratio.



**Fig. 14** Pipe Bend crack at intrados external.



# Fig. 15 SIF variation w.r.t. bend radius ratio for different values of k and a/t=0.75.

















## Case IV: Axial part through crack at intrados internal

In the present case, pipe-bend dimensions, crack length, loading, boundary conditions, material properties and meshing remains the same as in the previous cases. The geometry of the pipe bend along with the crack and boundary conditions is shown in Fig. 19. The dimensions of interaction integral contour are taken same as case-III. A regular mesh of 5x100x18 nodes is used for all simulations. The variation of SIF with pipe bend radius ratio ( $R^*$ ) is shown in Fig. 20 for four different values of cross-sectional radius ratio (k). The relative crack depth (a/t) is kept constant, and is taken as 0.75. The inner radius of the pipe bend ( $R_1$ ) is kept constant for these simulations. The variation of  $K_1$  with  $R^*$  is shown in Figs. 20(a), 20(b), 20(c) and 20(d) for k equal to 1.25, 1.375, 1.5 and 1.625, respectively. From these figures, it can be seen that the SIF values marginally decreases with the increase in bend radius ratio from 5.0 to 8.75.

Fig. 21 depicts the variation of stress intensity factor with  $R^*$  for four different values of k. The value of a/t is kept constant and is taken as 0.5. Figs. 21(a), 21(b), 21(c) and 21(d) present the variation of SIF with  $R^*$  for k equal to 1.25, 1.375, 1.5 and 1.625, respectively. From these results, it is concluded that the SIF values marginally decreases with the increase in  $R^*$  from 5.0 to 8.75. It is also observed from these plots that the high value of relative crack depth results in more SIF values if rest of the pipe-bend parameters is kept constant. The variation of SIF with radius ratio (k) is shown in Fig. 22 for four different values of bend radius ratio ( $R^*$ ). The relative crack depth (a/t) is again kept constant and is taken as 0.75. The inner radius of the pipe bend  $(R_1)$  is kept same for these simulations. Figs. 22(a), 22(b), 22(c) and 22(d) show the variation of SIF with radius ratio for the bend radius ratio of 5.0, 6.25, 7.5 and 8.75, respectively. From these figures, it can be seen that the SIF values decreases with the increase in radius ratio from 1.25 to 1.625.Fig. 23 depicts the variation of stress intensity factor with respect to radius ratio for four different values of bend radius. The relative crack depth (a/t) and inner radius of the pipe bend  $(R_1)$  are kept constant for these simulations. Figs. 23(a), 23(b), 23(c) and 23(d) present the variation of SIF with radius ratio for the bend radius ratio of 5.00, 6.25, 7.5 and 8.75, respectively. These plots show that the SIF values decreases with the increase in bend radius ratio from 1.25 to 1.625. On the basis of these simulations, it is observed that the numerical results obtained by XFEM are found quite close to the results obtained by FEM (Ansys). These results show that the large relative crack depth results in large SIF values if rest of the pipe bend parameters is kept constant. The influence of cross sectional radius ratio is found quite large in comparison to bend radius ratio. It is also noticed that the values of SIF are higher for an externally exposed crack in comparison to an internally exposed crack.

The present analysis also reveals that XFEM easily simulate 3-D cracks in pipe-bend in comparison to FEM as it does not require conformal meshing. In XFEM, these crack problems are analyzed by a regular mesh only. Moreover, the crack growth study in pipe-bend can be easily performed by XFEM as compared to FEM since FEM will not only require a conformal mesh but also will require re-meshing.



**Fig. 19** Pipe Bend crack at intrados internal.







**Fig.21** SIF variation w.r.t. bend radius ratio for different values of *k* and a/t=0.50.





• -XFEM

0

- FEM(Ans

1.6

- XFEM FEM(Ans

1.6

1.7

1.7

Fig. 22 SIF variation w.r.t. radius ratio for different values of  $R^*$  and a/t=0.75.





# 4 CONCLUSIONS

In this work, a detailed analysis of an axial semi-elliptical crack in pipe bend is performed by XFEM. The stress intensity factors for a semi-elliptical part through thickness crack placed on inner and outer surface of the pipe bend are evaluated by XFEM under internal pressure. In XFEM, a crack is easily modeled by enrichment functions without its physical presence using a regular mesh. The values of SIFs are obtained at four different locations of the pipe-bend for a wide range of geometric parameters like bend radius ratio, cross sectional radius ratio and relative crack length. The SIFs are extracted from XFEM solution by domain based interaction integral approach.

These results show that the location and type of the crack in pipe-bend has a significant effect on SIFs. The presence of crack at extrados location significantly affects the stress intensity factors as compared to intrados location. It is also observed that the large relative crack depth results in large SIFs if the rest of parameters are kept constant. The influence of cross sectional radius ratio is found more in comparison to bend radius ratio. It is also seen that the values of SIF is higher for an externally exposed crack in comparison to internally exposed crack. In all these simulations, the results obtained by XFEM were found in good agreement with those obtained by FEM. This work can be further extended for the simulation of circumferential cracks in pipe-bend. Moreover, the crack growth study in pipe-bend can also be easily performed by XFEM as it does not require re-meshing.

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