



Investigation of Stress in the Perforated Plate with the Presence of Edge Crack

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Abstract

The phenomenon of failure in objects is one of the major issues that human beings have been facing for a long time, and because of advances in technology in the present age, this issue is more important than in the past. All engineering materials, on the other hand, have tiny cracks from which failure begins. Therefore, estimating the residual life of thin plates made from these materials and used in space and offshore structures requires knowledge of the stress distribution due to cracking in these components. Because of the singularity of the crack tip due to large stresses, the presence of a relatively small crack can lead to a hazardous situation. Therefore, this area should be given more attention. In this research, using the Inglis formula and considering the correction coefficient of the compensatory free surface, the value of the stress coefficient for edge crack is obtained. Then, by replacing the new stress coefficient in Westergaard formula, we calculate the stress field of Mode I (Opening mode) and Mode II (sliding mode) in the perforated plane containing the edge crack. Finally, we examine the effects of various parameters such as loading angle, crack length and hole radius on the values obtained for stress in both modes by plotting.

Keywords: failure; edge crack; Inglis formula; stress field

1. Introduction

Thin sheets are widely used in offshore and space structures. Usually, defects in engineering structures or structures are involuntarily or specifically designed for structures. One of these geometrical defects is to create holes in the plates. Different types of holes are created in the plates for functional reasons such as weight loss of the structure and access to other system equipment. Plates containing holes when exposed to tensile or shear loading have caused severe local stress and reduced structural resistance, resulting in unexpected hazardous events such as the collapse of planes, spacecraft and ships, etc. over the past decades. This necessitates a more accurate understanding of the mechanics of failure in objects, and the problem of failure mechanics since the beginning of the present century has always been of interest to researchers and designers of mechanical components. Application of Complex Variable Functions to Plate Problems Based on Klovsov's Elasticity Theory was studied [1]. Investigation of inherent defects in materials, and stress fields in elliptical perforated plates by Inglis was showed [2]. Application of the complex mapping method in the problem of circular hole with edge crack and also circular hole with a pair of symmetric edge crack in unlimited plate under uniform stress was studied by Bowie [3]. The expression of the mixed-mode method in two-dimensional elastic theory by Muskhelishvili and his use of coherent mapping to analytically solve the stress problem around the hole in two-dimensional elastic bodies [4, 5]. the stress concentration around holes of different shapes (triangular-rectangular, etc.) for isotropic and non-isotropic planes using the Savin Measured to calculate [6]. the definition of special function

expansions, the application of the improved cumulative boundary method, and the Muskhelishvili mixing variable method, and computation for Newman's elliptical and circular bifurcations was used [7]. Numerical calculation of stress intensity coefficients using integral equations for circular perforations having an edge elasticity in infinite elastic material under uniform loading and also analysis of stress intensity factors under conditions where there is loading on the hole and crack by Tweed and Ruck was studied [8, 9]. tensile problems in a plate containing an elliptic hole through which two symmetric edge cracks are branched, and numerical results for mid-length cracks by Neyestani and Isida [10, 11] was Analyzed. the Semi-Unlimited and Isotropic Plate Problem of a Triangular Slit, from which a crack originated from a distal tensile load, and the use of the Muskhelishvili method and the congruent mapping as a sum of fractional expressions to provide an exact solution for the stress distribution by Hasseeb [12] And Ida was Investigated. Use of Green's Function and Muskhelishvili's Method to Calculate Stress Coefficients of First and Second Mode Stress Intensity Coefficients for an Unlimited Isotropic Plate Containing a Circular Hole with an Edge Crack by Shivakumar and Foreman was studied [13]. The stress distribution and intensity around triangular holes by Theocariz and Petro was Investigated [14]. Unlimited Isotropic Plate Problem Containing Cracked Circular Cracks and Crack Displacement Rate Using Doddell Plastic Area Model for Small Cracks by Bostrom was showed [15]. the Buckner principle to compute the stress intensity factor of the first and second mode infinite isotropic planes containing the cracked circular hole, also investigate the problem of having one crack and two cracks of different lengths and

two cracks of arbitrary angle originating from the hole by Lin et al was Applied [16]. the fractional numerical method, in calculating the stress intensity coefficients in cracks originating from circular and square holes in an isotropic boundless plate by Yan[17, 18] was Used. a convex mapping to investigate the problem of a perforated edge crack in a linear elastic plane under in-plane elongation and using this mapping to move the outside of the cracked semicircle into a single circle by Abdulmola et al[19] was Introduced. the stress distribution adjacent to a hole and a crack for plane problems with crack originating from an elliptic hole Also use the mixed method to calculate the stress intensity coefficients of single and double edge cracks originating from an elliptic hole by Guo et al[20-24] was Obtained. the stress distribution problem around a polygonal hole in an infinite plane under uniform load over long distances using the Batista-modified Muskhelishvili variable method[25] was Solved. The stress distribution around polygon hole in infinite plane under dual axis loading and investigation of the effect of hole geometry and loading by Sharma[26] was Calculated. the mixed variable method, in the study of the problem of elliptical holes with two edge cracks with unequal length in the isotropic plate and under tensile load in the distance and obtaining analytical solutions for stress functions and stress intensity coefficient of the first mode by Liu and Duan[27-32] was Used. the fractional method, to calculate stress intensity coefficients on a rectangular plate with two equal cracks by Miao et al[33, 34] was studied. The First Mode Stress Intensity Coefficients in Single and Double Circular Plate Finite Circular Plates Using New Weight Function by Kim and Hill[35] was Presented. The fracture behavior due to an elliptical hole with two asymmetric cracks in a one-dimensional quasi-crystal hexagon with piezoelectric effect by Yang et al[36] was Investigated. In this research, by combining the Inglis formula with the Westergaard method and considering the correction factor of the compensatory free surface, we introduce a new method for calculating the stress field of a perforated plate containing edge cracks for modes 1 and 2, Then the effects of crack length, loading angle and hole radius on the stress level for both modes are compared.

2. Problem Statement

Suppose an indefinite plate containing crack length $2a$ under uniform tension σ locates in the far distance (Fig. 1). For this problem, the simplest function of Westergaard is as follows[37].

$$z_{II} = \frac{\sigma z}{(z^2 - a^2)^{1/2}} \quad (1)$$

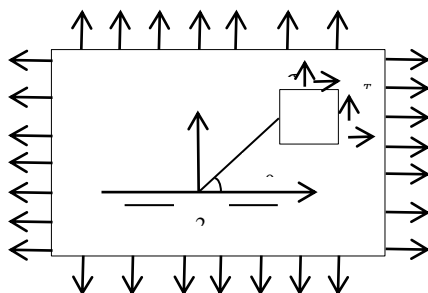


Fig.1. Infinite plate containing cracks of length $2a$ under uniform stresses σ at infinity

If in Fig. 1, instead of normal tension, σ applies uniform shear tension, the simplest Westergaard function is The following is the case[37]:

$$z_{II} = -\frac{itz}{(z^2 - a^2)^{1/2}} \quad (2)$$

2.1. A) Mode I (Opening mode)

Consider Equation (1). If we transfer the source of the coordinates to the crack tip, then $\xi = z - a$ Then the equation for points near the crack tip ($\xi \ll a$) can be calculated as follows:

$$\begin{aligned} z_I &= \frac{\sigma(\xi + a)}{[\xi(\xi + 2a)]^{1/2}} = \frac{\sigma(\xi + a)}{\xi^{1/2} \cdot (\xi + 2a)^{1/2}} \\ &= \frac{\sigma(\xi + a)}{\xi^{1/2} \left(2a \left(1 + \frac{\xi}{2a}\right)\right)^{1/2}} = \\ &= \frac{\sigma(\xi + a)}{(2a\xi)^{1/2} \left(1 + \frac{\xi}{2a}\right)^{1/2}} \\ &= \frac{\sigma(\xi + a)}{(2a\xi)^{1/2}} \left[1 - \frac{1}{2} \left(\frac{\xi}{2a}\right) + \frac{3}{8} \left(\frac{\xi}{2a}\right)^2 - \frac{15}{24} \left(\frac{\xi}{2a}\right)^3 + \dots\right] \end{aligned}$$

But around the crack tip which is $\xi = z - a$ (values $|\xi|$ are very small, so assuming $k_1 = \sigma\sqrt{\pi a}$ can be written as:

$$\begin{aligned} z_I &= \frac{\sigma a}{(2a\xi)^{1/2}} = \frac{\sigma\sqrt{a}}{(2\xi)^{1/2}} = \frac{\sigma\sqrt{\pi a}}{\sqrt{2\pi\xi}} \\ &= \frac{k_1}{\sqrt{2\pi\xi}} \quad (3) \end{aligned}$$

Now by considering the polar coordinates r, θ , and $\xi = re^{i\theta}$ apply the law of motion.

$$\sigma_x = \frac{k_1}{\sqrt{2\pi r}} \cos \frac{\theta}{2} \left(1 - \sin \frac{\theta}{2} \sin \frac{3\theta}{2}\right) \quad (4)$$

$$\sigma_y = \frac{k_1}{\sqrt{2\pi r}} \cos \frac{\theta}{2} \left(1 + \sin \frac{\theta}{2} \sin \frac{3\theta}{2}\right) \quad (5)$$

$$T_{xy} = \frac{k_1}{\sqrt{2\pi r}} \cos \frac{\theta}{2} \left(\sin \frac{\theta}{2} \cos \frac{3\theta}{2}\right) \quad (6)$$

B) II (sliding mode)

By applying the relation (2) and the process used for the first mode, different tension states can be calculate for the crack with the length of $2a$ in an unlimited sheet under uniform shear tension in the plate (τ) as follow:

$$\sigma_x = \frac{k_{II}}{\sqrt{2\pi r}} \sin \frac{\theta}{2} \left(2 + \cos \frac{\theta}{2} \cos \frac{3\theta}{2} \right) \quad (7)$$

$$\sigma_y = \frac{k_{II}}{\sqrt{2\pi r}} \sin \frac{\theta}{2} \left(\cos \frac{\theta}{2} \cos \frac{3\theta}{2} \right) \quad (8)$$

$$T_{xy} = \frac{k_{II}}{\sqrt{2\pi r}} \cos \frac{\theta}{2} \left(1 - \sin \frac{\theta}{2} \sin \frac{3\theta}{2} \right) \quad (9)$$

Where (k_{II}) the tension coefficient of mode II is for the following figure.

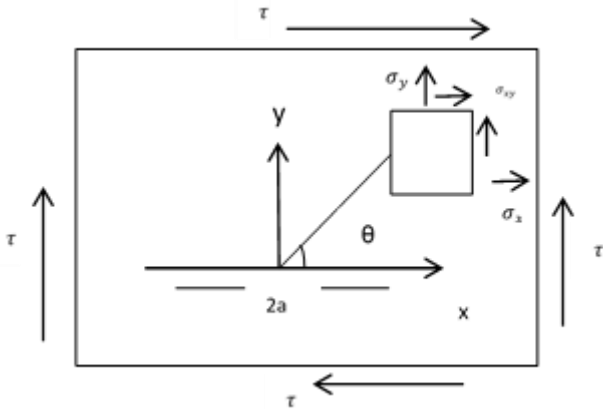


Fig.2. Infinity plate containing cracks of length 2a under uniform shear stresses τ at infinity

$$(k_{II} = \tau\sqrt{\pi a}) \quad (10)$$

On the other hand, taking into account the British equation [2], the tension distribution in the perforated plate containing the elliptic hole For Fig. 3, can be obtained according to the following formula.

$$\sigma_A = \sigma \left(1 + \frac{2a}{b} \right) \quad (11)$$

This relationship gives us the amount of tension on a plate with elliptic holes 2a and 2b. Also according to [38]. The stress intensity coefficient for a fracture crack is obtained from relation (11) that its σ is the applied stress and a_0 is the length of the fracture crack.

$$k = 1.12\sigma_A\sqrt{\pi a_0} \quad (12)$$

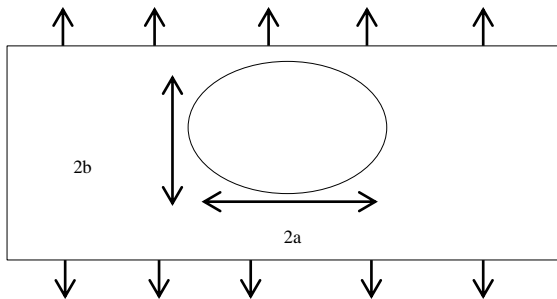


Fig.3. The plate has an elliptical hole of diameter 2 a, 2b under uniform stresses σ at infinity

Now, if we consider the simultaneous effects of holes and cracks, we can integrate the relationships (11) and (12) for obtaining the following relation:

$$k^* = 1.12\sigma \left(1 + \frac{2a}{b} \right) \sqrt{\pi a_0} \quad (13)$$

Of course, the relation (13) for the time ($a_0 \ll R$) means that the dimensions of the crack are very low compared to the cavity radius.

In the particular case where $a = b = R$. With regard to the relation (12) we will have:

$$k_1^* = 3.36\sigma\sqrt{\pi a_0} \quad (14)$$

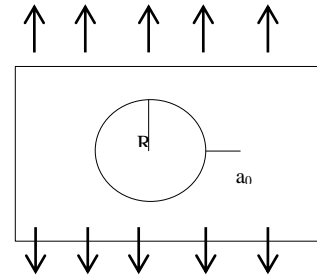
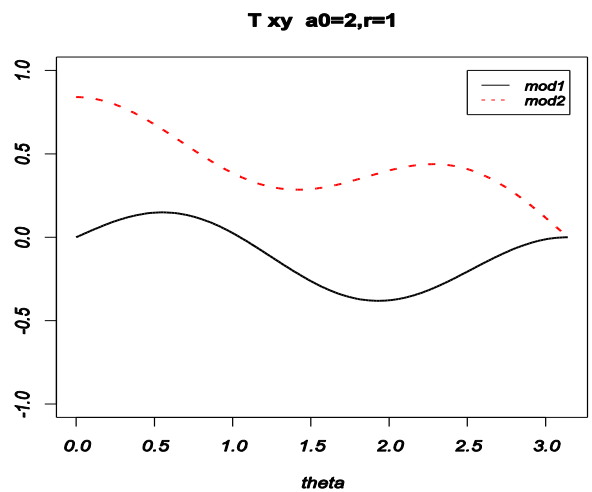


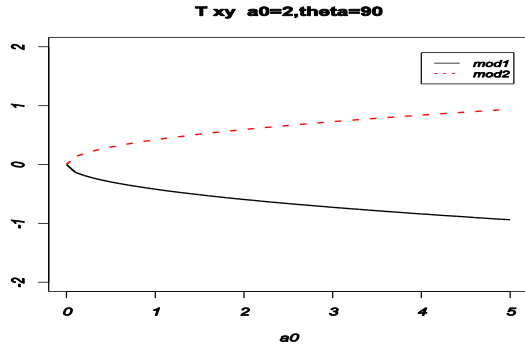
Fig.4. Short cracks of length a_0 , caused by an elliptical hole in an infinite sheet subjected to stress σ perpendicular to the principal axis of the hole

By substituting k_1^* for k_1 in relation (4), (5) and (6) the stress field of mode II can be calculated on perforated plates containing the edge crack for Fig.4. Also, if instead of the normal stress σ , set the value of τ in the relation k_1^* . then put the obtained value in the equations (7), (8) and (9). The Mode II stress value for the perforated plate containing the crack is obtained. Now to investigate the effects of various parameters such as loading angle, crack length and hole radius on mode stress values 1 and 2 using R coding software version 3.6.3 Draw and analyze the results.

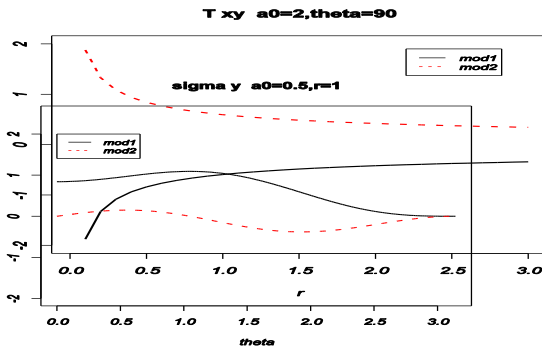
3. Results and discussion



(a)



(b)

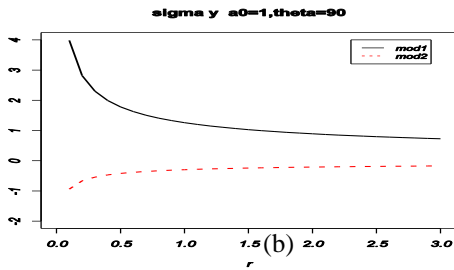


(c)

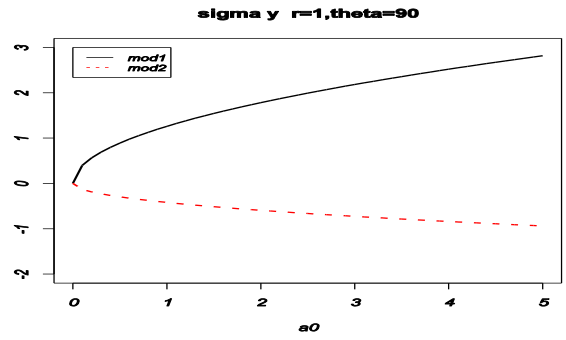
Fig.5.a,b,c respectively The effect of θ , a_0 , and r on T_{xy} in modes I and II

Fig.5 compare the rate of change of T_{xy} in modes I and II and in cases where a_0 , r , and θ change. As we can see in the figures above, with increasing angle θ , however, the value of T_{xy} in mode II is greater than that in mode I But these changes are oscillating, and increasing or decreasing the angle does not increase or decrease the amount of T_{xy} at all. On the other hand, increasing a_0 in mode I reduces T_{xy} but in mode II increases T_{xy} . In general, as the value of a_0 increases, the value of T_{xy} diverges in two modes. Also, increasing r in mode I increases T_{xy} but in mode II decreases T_{xy} . And in general, as r increases, the difference in T_{xy} value between the two modes decreases.

(a)



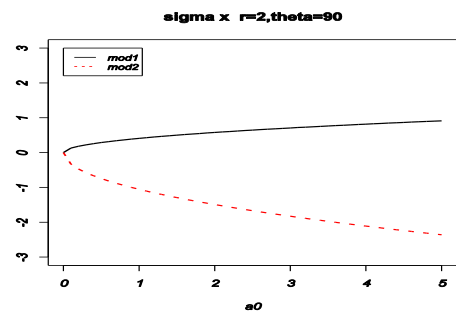
(b)



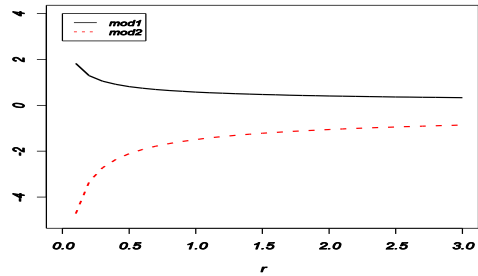
(c)

Fig.6. The effect of r , a_0 and θ on σ_y in modes I and II
 Fig.6.compare the rate of change of σ_y in modes I and II and in cases where a_0 , r , and θ change. As we can see in the above figures, by increasing r , the value of σ_y decreases in mode I but increases in mode II. On the other hand, the higher the value of r , the closer the value of σ_y becomes in two modes. But with increasing a_0 , the value of σ_y increases in case I and decreases in mode II but, with increasing angle θ , however, the value of σ_y in mode I is greater than that in mode II. And then, with a proper increase in angle θ , the two values converge.

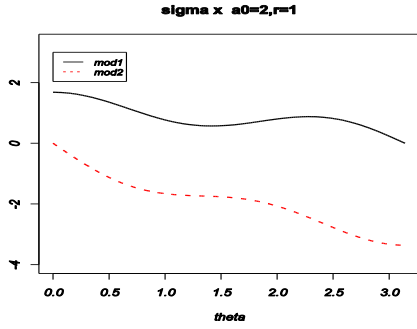
(a)



(b)



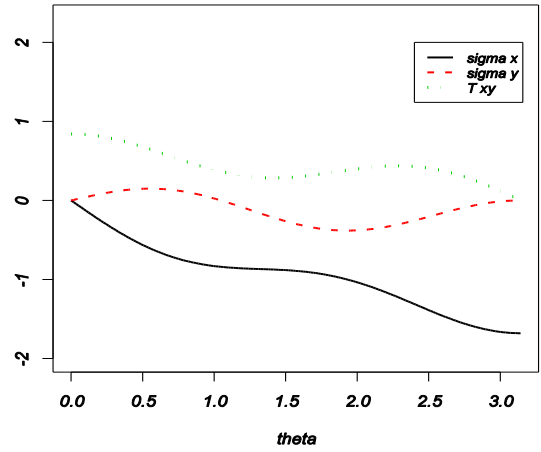
(b)



(c)

Fig.7. The effect of r , a_0 and θ on σ_x in modes I and II
 Fig.7 show the rate of change of σ_x in states where a_0 , r , and θ change. As shown in the diagrams. Increasing r in mode I decreases σ_x but in mode II increases σ_x . But in the case of increasing a_0 , the opposite is true, ie in mode I the value of σ_x increases but in mode II it decreases. Also, as the angle. Increases, the difference in the value of σ_x between the modes increases.

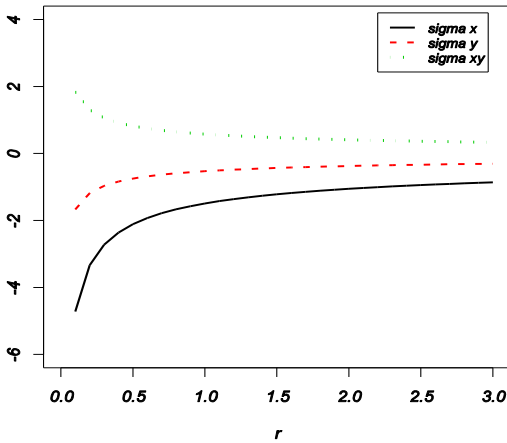
mod2 a0=0.5,r=1



(c)

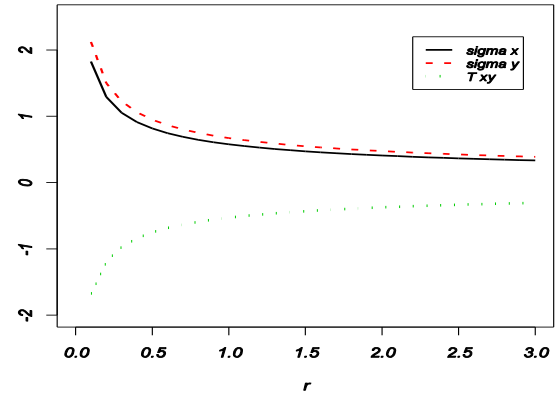
Fig.8. The effect of r , a_0 and θ on σ_x , σ_y and T_{xy} in mode II

mod2 a0=1,theta=90



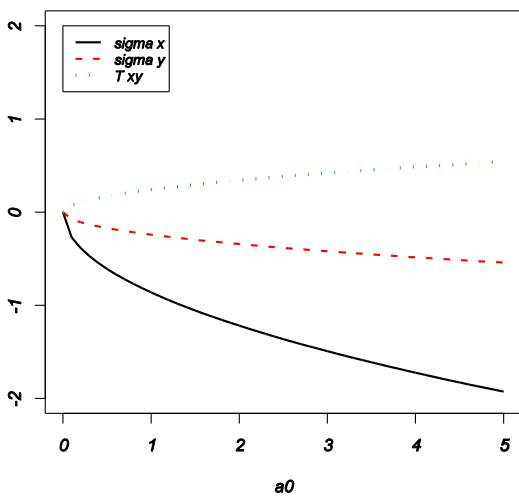
(a)

mod1 a0=1,theta=90



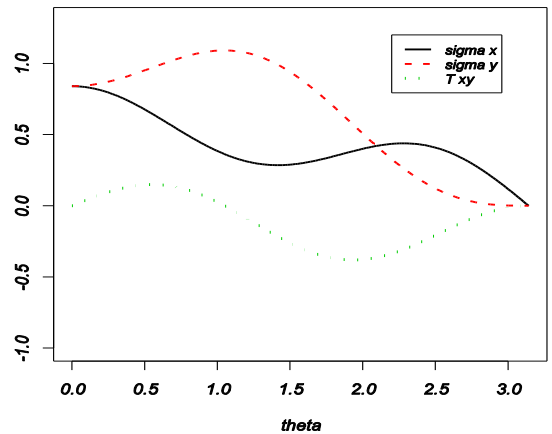
(a)

mod2 r=3,theta90

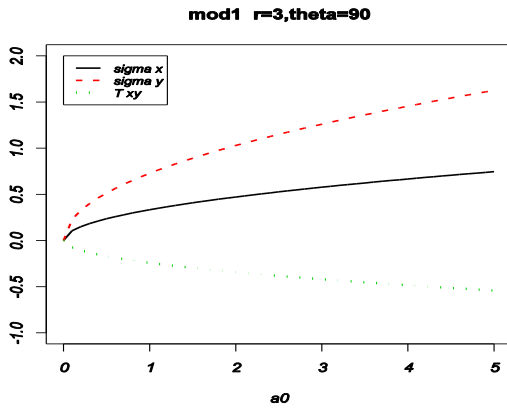


(b)

mod1 a0=0.5,r=1



(b)



(c)

Fig.9. The effect of r, a_0 and θ on σ_x, σ_y and T_{xy} in mode I
 Fig.8 Comparison diagrams of σ_x, σ_y and T_{xy} values in case II and Fig.9 show a comparison chart of the values mentioned in Mode I

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4. Conclusion

As the value of r increases, σ_x has a decreasing trend in mode I but an increasing trend in mode II. T_{xy} also increases in mode I but decreases in mode II. And σ_y increases in mode II but decreases in mode I. Increasing the value of a_0 in mode I decreases T_{xy} and increases σ_x and σ_y . While in Mode II the trend is different, ie T_{xy} increases but σ_x and σ_y decrease. By increasing θ in both modes in the considered range, we do not see ever decreasing or ascending behavior of any of the values σ_x, σ_y and T_{xy} . Rather, the behavior of the relevant graphs in different intervals is sometimes ascending and sometimes descending. And it can not be said at all that increasing θ in both modes increases or decreases the values of σ_x, σ_y and T_{xy} .

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